# Duality structures on tensor categories coming from vertex operator algebras 

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## Overview

(1) Vertex (operator) algebras and commutative algebras
(2) Tensor products
(3) Duality structures

## Definition: Vertex algebra (VA)

## Data:

- $\mathbb{C}$ vector space $V$
- field map
$Y_{z}: V \otimes V \rightarrow V(z)$
- vacuum vector $\Omega \in V$.

Axioms:

- vacuum axiom:

$$
\begin{aligned}
& Y_{z}(\Omega \otimes-)=\mathrm{id}_{V} \text { and } \\
& Y_{z}(A \otimes \Omega)=A+z V \llbracket z \rrbracket, \forall A \in V
\end{aligned}
$$

- locality: $\forall A, B \in V, \exists n \in \mathbb{N}$

$$
(z-w)^{n}\left[Y_{z}(A \otimes-), Y_{w}(B \otimes-)\right]=0
$$

## Consequence/Proposition

- There exists a translation operator $T: V \rightarrow V$, such that $T \Omega=0$ and $\left[T, Y_{z}\right]=\partial_{z} Y_{z}$
- For all $A, B, C \in V$

$$
\begin{array}{r}
Y_{z}\left(A \otimes Y_{w}(B \otimes C)\right) \in V(z)(w) \\
Y_{w}\left(B \otimes Y_{z}(A \otimes C)\right) \in V(w)(z) \\
Y_{w}\left(Y_{z-w}(A \otimes B) \otimes C\right) \in V(w)(z-w)
\end{array}
$$

are expansions of same element in $V \llbracket z, w \rrbracket\left[z^{-1}, w^{-1},(z-w)^{-1}\right]$.

## Example: Heisenberg vertex algebra

- Heisenberg Lie algebra: $\mathfrak{h}=\bigoplus_{n \in \mathbb{Z}} \mathbb{C} a_{n} \oplus \mathbb{C} 1$ with 1 central and relations $\left[a_{m}, a_{n}\right]=m \delta_{m, n} \mathbf{1}$.
- Fock space (h Verma module), $F_{\lambda}=\mathbb{C}\left[a_{-n}: n \geq 1\right]|\lambda\rangle$, where $a_{0}=\lambda \mathrm{id}, a_{n} / n=\partial_{a_{-n}}, n \geq 1$.
- The assignments and recursions
$|0\rangle \otimes-\mapsto \mathrm{id}_{F_{0}}, \quad a_{-1}|0\rangle \otimes-\mapsto \sum_{i \in \mathbb{Z}} a_{i} z^{-i-1}=a(z)$,
$a_{-n}|0\rangle \otimes-\mapsto \frac{\partial^{n-1}}{(n-1)!} a(z) \otimes-$,
$a_{-n} v \otimes-\mapsto\left(\frac{\frac{2}{}_{n-1}^{(n-1)!}}{(n)}(z)\right)_{r} Y_{z}(v \otimes-)+Y_{z}\left(v \otimes\left(\frac{\partial^{n-1}}{(n-1)!} a(z)\right)_{\mathrm{p}}-\right)$.
define a vertex algebra on $F_{0}$ with vacuum vector $|0\rangle$.


## Definition: Vertex operator algebra (VOA)/ conformal vertex algebra

A vertex operator algebra is a vertex algebra $(V, Y, \Omega, T)$ admiting a conformal vector $\omega \in V$ satsifying

- Virasoro algebra relations: $Y_{z}(\omega \otimes-)=\sum_{n \in \mathbb{Z}} L_{n} z^{-n-2}$

$$
\left[L_{m}, L_{n}\right]=(m-n) L_{m+n}+\frac{c}{12}\left(m^{3}-m\right) \delta_{m+n, 0},
$$

with $c \in \mathbb{C}$, central charge.

- Non-negative integral conformal grading:

$$
V=\bigoplus_{n=0}^{\infty} V_{n}, V_{n}=\left\{v \in V \mid\left(L_{0}-n\right) v=0\right\} .
$$

- Conformal derivation: $L_{-1}=T$.


## Example: Heisenberg vertex operator algebra

Consider the vertex algebra $\left(F_{0}, Y\right)$ from before.

- For $\rho \in \mathbb{C}$, the vector
$\omega_{\rho}=\left(\frac{1}{2} a_{-1}^{2}+\rho a_{-2}\right)|0\rangle \mapsto \frac{1}{2}\left(a(z)_{\mathrm{r}} a(z)+a(Z) a(z)_{\mathrm{p}}\right)+\rho \partial a(z)$
is conformal.
- The central charge of the resulting Virasoro algebra is $c_{\rho}=1-12 \rho^{2}$.
- The conformal grading on $F_{0}$ assigns grade 0 to $|0\rangle$ and grade $-n$ to $a_{n}$.


## Recap

- Vertex algebras are essentially associative commutative unital $\mathbb{C}$-algebras with a derivation.
- Fields $Y_{z}: V \otimes V \rightarrow V(z)$ are essentially an action of $V$ on itself.
- Where you have commutativie algebras, you have modules!


## Definition: Vertex algebra module

Let $(V, \Omega, T, Y)$ be a vertex algebra. A $V$-module is a pair $\left(M, Y^{M}\right): M$ a vector space and $Y_{z}^{M}: V \otimes M \rightarrow M(z)$ a $V$-action, that is,

- $Y_{z}^{M}(\Omega \otimes-)=\operatorname{id}_{M}$
- For all $A, B \in V$ and $C \in M$ the expansions
$Y_{z}^{M}\left(A \otimes Y_{w}^{M}(B \otimes C)\right) \in M(z)(w)$
$Y_{w}^{M}\left(B \otimes Y_{z}^{M}(A \otimes C)\right) \in M(w)(z)$
$Y_{w}^{M}\left(Y_{z-w}(A \otimes B) \otimes C\right) \in M(w)(z-w)$
can be identified in $M[[z, w]]\left[z^{-1}, w^{-1},(z-w)^{-1}\right]$.
Many additional assumtions can be added to the above definition. E.g. bounded conformal weights, finite weight spaces, semi simplicity, etc.


## Example: Heisenberg modules

- Recall the Heisenberg algebra $\mathfrak{h}$ and Fock spaces $F_{\lambda}$ from before.
- For $\lambda \in \mathbb{C}$, the Fock space $F_{\lambda}$ is a $\left(F_{0}, Y\right)$-module with the action defined by the same formula, e.g. $Y_{z}^{F_{\lambda}}\left(a_{-1}|0\rangle \otimes|\lambda\rangle\right)=a(z)|\lambda\rangle$.


## Duals of modules

## Definition: the dual of a module

Let $(V, Y, \Omega, \omega)$ be a vertex operator algebra and $\left(M, Y^{M}\right), M=\bigoplus_{n} M_{n}$ a module. Then $\left(M^{\prime}, Y^{M^{\prime}}\right)$

- $M^{\prime}=\bigoplus_{n} \operatorname{Hom}_{\mathbb{C}}\left(M_{n}, \mathbb{C}\right)$,
- $\left\langle Y_{z}^{M^{\prime}}(v \otimes \mu), m\right\rangle=\left\langle\mu, Y_{z^{-1}}^{M}\left(e^{z L_{1}}\left(-z^{2}\right)^{L_{0}} v \otimes m\right)\right\rangle, v \in V, m \in M, \mu \in M^{\prime}$. is again a $(V, Y)$-module.

Heisenberg example:
For the choice of conformal vector $\omega_{\rho}$, we have $F_{\lambda}^{\prime} \cong F_{2 \rho-\lambda}$

## Motivating tensor products

- In quantum field theory all information is encoded in $n$-point correlation functions.
- In conformal quantum field theory (CFT) these correlation functions are $V$-multilinear functions.
- So we need to understand multilinear algebra for vertex algebras.
- This is also a natural question for commutative algebras.


## Definition: Intertwining operator, $V$-bilinear maps

Let $(V, \Omega, Y, \omega)$ be a vertex operator algebra and
$\left(M_{1}, Y^{M_{1}}\right),\left(M_{2}, Y^{M_{2}}\right),\left(M_{3}, Y^{M_{3}}\right)$ be $V$-modules. An intertwining operator of type $\left(\begin{array}{c}M_{1}, M_{2}\end{array}\right)$ is a map $\mathcal{Y}_{x}: M_{1} \otimes M_{2} \rightarrow M_{3}\{x\}$ such that for all $m_{i} \in M_{i}$

- $\mathcal{Y}_{x}\left(m_{1} \otimes m_{2}\right)$ truncates below.
- $\mathcal{Y}_{x}\left(L_{-1} m_{1} \otimes m_{2}\right)=\partial_{x} \mathcal{Y}_{x}\left(m_{1} \otimes m_{2}\right)$.
- The expansions $Y_{z}^{M_{3}}\left(A \otimes \mathcal{Y}_{x}\left(m_{1} \otimes m_{2}\right)\right) \sim \mathcal{Y}_{x}\left(Y_{z-x}^{M_{1}}\left(A \otimes m_{1}\right) \otimes m_{2}\right) \sim \mathcal{Y}_{x}\left(m_{1} \otimes Y_{z}^{M_{2}}\left(A \otimes m_{2}\right)\right)$ can be identified.

Observations:

- The field map $Y$ is an intertwining operator of type $\binom{V}{V, V}$.
- The action $Y^{M}$ is an intertwining operator of type $\left(\begin{array}{l}V, M\end{array}\right)$.
- Intertwining operators are $V$-bilinear maps. All intertwining operators of a given type form a vector space. The field map $Y$ and the action $Y^{M}$ have a distinguished normalisation due to $Y_{z}(\Omega \otimes-)=\mathrm{id}$.


## Example: Heisenberg intertwining operators

Recall the Heisenberg Fock spaces $F_{\mu}, \mu \in \mathbb{C}$.
Then $\operatorname{dim}\binom{F_{\rho}}{F_{\mu}, F_{\nu}}=\delta_{\rho, \mu+\nu}$ for all $\rho, \mu, \nu \in \mathbb{C}$.
$\binom{F_{\mu}+\nu}{F_{\mu}, F_{\nu}}$ is spanned by

$$
\begin{array}{r}
\mathcal{Y}_{x}^{F_{\mu}, F_{\nu}}(p|\mu\rangle \otimes q|\nu\rangle)=x^{\mu \nu} S_{\mu} \prod_{m \geq 1} \exp \left(\mu \frac{a_{-m}}{m} x^{m}\right) Y_{x}^{F_{\nu}}(p|0\rangle \otimes-) \\
\cdot \prod_{m \geq 1} \exp \left(-\mu \frac{a_{m}}{m} x^{-m}\right) q|\nu\rangle
\end{array}
$$

where $S_{\mu}:|\nu\rangle \mapsto|\mu+\nu\rangle$ is the shift operator.

Tensor products pull multilinear algebra back to linear algebra!

## Definition: Fusion product aka vertex algebra tensor product

Let $(V, \Omega, Y, \omega)$ be a vertex operator algebra and $\left(M_{1}, Y^{M_{1}}\right),\left(M_{2}, Y^{M_{2}}\right)$ be $V$-modules. A fusion product is a triple ( $M_{1} \boxtimes M_{2}, Y^{M_{1} \boxtimes M_{2}}, \mathcal{Y}^{M_{1}, M_{2}}$ ), where ( $M_{1} \boxtimes M_{2}, Y^{M_{1} \boxtimes M_{2}}$ ) is a $V$-module and $\mathcal{Y}^{M_{1}, M_{2}}$ is an intertwining operator of type $\binom{M_{1} \boxtimes M_{2}}{M_{1}, M_{2}}$ such that the following universal property holds: For every $V$-module ( $X, Y^{X}$ ) and intertwining operator $\mathcal{Y}^{X}$ of type $\left(\begin{array}{c}M_{1}, M_{2}\end{array}\right)$

In contrast to linear algebra (or ring theory) constructing $M_{1} \boxtimes M_{2}$ and decomposing into a direct sum of indecomposable modules is extremely hard.

Well chosen categories of modules are tensor categories with respect to $\boxtimes$ with the following structures. [Huang-Lepowsky-Zhang]

- For module homomorphimsms $f: X \rightarrow Z, g: U \rightarrow W$, the morphism $f \boxtimes g$ is uniquely characterised by $(f \boxtimes g) \circ \mathcal{Y}^{X \boxtimes U}=\mathcal{Y}^{Z \boxtimes W} \circ(f \otimes g)$
- $V$ is the tensor identity and the unit isomorphisms are uniquely characterised by
$\ell_{M}\left(\mathcal{Y}_{z}^{V, M}(a \otimes m)\right)=Y_{z}^{M}(a \otimes m)$ and
$r_{M}\left(\mathcal{Y}^{M, V}(m \otimes a)\right)=e^{z L_{-1}} Y_{-z}^{M}(a \otimes m)$.
- associativity isomorphisms (hardest part!)

$$
\begin{aligned}
& A_{M_{1}, M_{2}, M_{3}}\left(\mathcal{Y}_{x_{1}}^{M_{1}, M_{2} \boxtimes M_{3}}\left(m_{1} \otimes \mathcal{Y}_{x_{2}}^{M_{2}, M_{3}}\left(m_{2} \otimes m_{3}\right)\right)\right)= \\
& \mathcal{Y}_{x_{2}}^{M_{1} \boxtimes M_{2}, M_{3}}\left(\mathcal{Y}_{x_{1}-x_{2}}^{M_{1}, M_{2}}\left(m_{1} \otimes m_{2}\right) \otimes m_{3}\right)
\end{aligned}
$$

All analytic details hidden.

- Braiding isomorphisms uniquely characterised by

$$
c_{M_{1}, M_{2}}\left(\mathcal{Y}_{x}^{M_{1}, M_{2}}\left(m_{1} \otimes m_{2}\right)\right)=e^{x L_{-1}} \mathcal{Y}_{e^{i \pi} x}^{M_{2}, M_{1}}\left(m_{2} \otimes m_{1}\right)
$$

- If the vertex algebra $V$ is conformal (a vertex operator algebra) and the modules are chosen to be compatible with this conformal structure, then there is also a twist $\theta_{M}=\left.e^{2 \pi i L_{0}}\right|_{M}$, which satisfies the balancing equation
$\theta_{M_{1} \boxtimes M_{2}}=c_{M_{1}, M_{2}} \circ c_{M_{2}, M_{1}} \circ\left(\theta_{M_{1}} \boxtimes \theta_{M_{2}}\right)$
- Tensor categories of vertex operator algebra modules depend only very weakly on the conformal structure. Only the twist and taking duals depend on the conformal structure.

Theorem [Huang, Moore-Seiberg]: The Verlinde Conjecture
Let $(V, Y, \Omega, \omega)$ be a vertex operator algebra and Adm $V$ be the category of admissible $V$-modules. If
(1) $\operatorname{dim} V_{0}=1, \operatorname{dim} V_{-n}=0, \operatorname{dim} V_{n}<\infty, n \in \mathbb{N}$,
(2) $V$ is simple as a module over itself,
(3) $V \cong V^{*}$, self-dual,
(4) $\operatorname{dim} V / c_{2}(V)<\infty$, (a technical finiteness condition)
(5) $\operatorname{Adm}(V)$ is semisimple,
then Adm $V$ is a modular tensor category. Further the action of the modular group on the category (which determines Verlinde's formula) is equal (after a renormalisation) to the action of the modular group on module characters.

## Recap

- Vertex algebras are almost commutative unital algebras with derivations.
- The conformal vector is a choice/structure: there can be 0,1 or many.
- Vertex algebras admit modules. "Good choices" of module categories admit a tensor (aka fusion) product.
- With the exception of associators, the tensor structure morphisms follow from easy constructions.


## Beyond modular tensor categories

The Verlinde conjecture is about vertex operator algebras that are maximally nice. There are many deviations from niceness.

- The category of modules need not be finite (the Heisenberg example is not finite).
- The category of modules need not be semisimple.
- The tensor product need not be left exact (there can be non-flat objects). [Gaberdiel-Runkel-SW]
- The first two deviations above can still be nice in the sense that they admit rigid duals. [Tsuchiya-SW,Allen-SW]
The third deviation is more fundamentally broken, as non-flat objects cannot have rigid duals.


## Duals of intertwining operators

## Proposition [HLZ]

Let $\mathcal{Y}$ be an intertwining operator of type $\left(\begin{array}{c}M_{1}, M_{2}\end{array}\right)$. Then

$$
\begin{aligned}
& M_{1} \otimes M_{3}^{\prime} \rightarrow M_{2}^{\prime}\{x\} \\
& m \otimes \nu \mapsto \nu\left(\mathcal{X}_{x^{-1}}\left(e^{z L_{1}}\left(-z^{2}\right)^{L_{0}} m \otimes-\right)\right)
\end{aligned}
$$

is an intertwining operator of type $\binom{M_{1}, M_{3}^{\prime}}{M_{2}^{\prime}}$.
Thus $\binom{M_{3}}{M_{1}, M_{2}} \cong\binom{M_{2}^{\prime}}{M_{1}, M_{3}^{\prime}}$.

## Remark

Intertwining operators of type $\binom{M_{3}}{M_{1}, M_{2}}$ should be thought of as Hom-spaces of the form $\operatorname{Hom}_{V}\left(M_{1} \boxtimes M_{2}, M_{3}\right)$.

## Defininition: Grothendieck Verdier categories

Let $\mathcal{C}$ be a monoidal (abelian linear) category.
(1) An object $K \in \mathcal{C}$ is called dualising, if for all $Y \in \mathcal{C}$ the functor $\operatorname{Hom}_{V}(-\otimes Y, K)$ is representable, that is,

$$
\operatorname{Hom}_{V}(X \otimes Y, K) \cong \operatorname{Hom}_{V}(X, G Y)
$$

and if the so defined contravariant functor $G: \mathcal{C} \rightarrow \mathcal{C}$ is an anti-equivalence.
(2) A monoidal category $\mathcal{C}$ together with a choice of dualising object $K \in \mathcal{C}$ is called a Grothendieck-Verdier or $*$-autonomous category.

Theorem [Allen-Lentner-Schweigert-SW]
Let $(V, Y)$ be a vertex operator algebra and let $\operatorname{Rep}(V)$ be choice of modules to which the HLZ tensor product theory applies which is in addition closed under taking duals. Then $V^{\prime}$ is a dualising object for $\operatorname{Rep}(V)$ and $\left(\operatorname{Rep}(V), V^{\prime}\right)$ is a Grothendieck-Verdier category (actually ribbon Grothendieck-Verdier).

## Why am I excited about this?

Grothendieck-Verdier categories have many appealing features

- The appear to be describe the natural duality structure of vertex operator algebra modules and the can accommodate all non-nice features mentioned previously.
- They admit two tensor products $X \otimes Y$ and $X \bullet Y=G\left(G^{-1} Y \otimes G^{-1} X\right)$, where $\otimes$ is right exact and $\bullet$ is left exact. In particular there are distributor morphisms
$\partial_{X, Y, Z}^{l}: X \otimes(Y \bullet Z) \rightarrow(X \otimes Y) \bullet Z, \quad \partial_{X, Y, Z}^{r}:(X \bullet Y) \otimes Z \rightarrow X \bullet(Y \otimes Z)$,
which need not be isomorphisms. [Fuchs-Schaumann-Schweigert-SW]
- $\otimes$ admits inner Homs (right adjoints) and • admits inner coHoms (left adjoints which allow the construction of algebras and coalgebras (even Frobenius algebras) in $\mathcal{C}$ (crucial for conformal field theory).

Thank you!

## For further reading I

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